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Equations (24) and (25) in Appendix B of [1] should respectively read as

$$\int_{V_\infty} -(\nabla G_1)G_2^* - \hat{\boldsymbol{r}} jk \frac{e^{jk(\boldsymbol{r}_1-\boldsymbol{r}_2)\cdot\hat{\boldsymbol{r}}}}{16\pi^2|\boldsymbol{r}|^2}\mathrm{dV} = -\frac{\boldsymbol{r}_{12}}{|\boldsymbol{r}_{12}|}\frac{\cos(k|\boldsymbol{r}_{12}|)}{8\pi}$$
$$-j\frac{2\boldsymbol{r}_1}{8\pi k^2}\left(\frac{\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^3} - \frac{k\cos(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^2}\right)$$
$$-j\frac{|\boldsymbol{r}_1|^2-|\boldsymbol{r}_2|^2}{8\pi k^2}\frac{\boldsymbol{r}_{12}}{|\boldsymbol{r}_{12}|^2}\left(\frac{k^2\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|}\right.$$
$$\left.-3\left(\frac{\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^3} - \frac{k\cos(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^2}\right)\right) \tag{1}$$

and

$$\int_{V_\infty} j(\nabla G_1)G_2^* - \hat{\boldsymbol{r}} k \frac{e^{jk(\boldsymbol{r}_1-\boldsymbol{r}_2)\cdot\hat{\boldsymbol{r}}}}{16\pi^2|\boldsymbol{r}|^2}\mathrm{dV} = j\frac{\boldsymbol{r}_{12}}{|\boldsymbol{r}_{12}|}\frac{\cos(k|\boldsymbol{r}_{12}|)}{8\pi}$$
$$-\frac{\boldsymbol{r}_{12}}{8\pi k^2}\left(\frac{\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^3} - \frac{k\cos(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^2}\right)$$
$$-\frac{\boldsymbol{r}_1+\boldsymbol{r}_2}{8\pi k^2}\left(\frac{\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^3} - \frac{k\cos(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^2}\right)$$
$$-\frac{|\boldsymbol{r}_1|^2-|\boldsymbol{r}_2|^2}{8\pi k^2}\frac{\boldsymbol{r}_{12}}{|\boldsymbol{r}_{12}|^2}\left(\frac{k^2\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|}\right.$$
$$\left.-3\left(\frac{\sin(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^3} - \frac{k\cos(k|\boldsymbol{r}_{12}|)}{|\boldsymbol{r}_{12}|^2}\right)\right) \tag{1a}$$
$$= j\frac{\boldsymbol{r}_{12}}{2}\mathrm{Re}\{G_{12}\} - \frac{1}{2k^2}\mathrm{Im}\{\nabla_1 G_{12}\}$$
$$-\frac{\boldsymbol{r}_1+\boldsymbol{r}_2}{2k^2}\mathrm{Im}\{\nabla_1 G_{12}\cdot\frac{\boldsymbol{r}_{12}}{|\boldsymbol{r}_{12}|^2}\}$$
$$+\frac{|\boldsymbol{r}_1|^2-|\boldsymbol{r}_2|^2}{2k^2|\boldsymbol{r}_{12}|^2}\mathrm{Im}\{\boldsymbol{r}_{12}k^2G_{12} + 3\nabla_1 G_{12}\}. \tag{1b}$$

All other results in [1] do not involve the coordinate-dependent terms (those with $\boldsymbol{r}_1+\boldsymbol{r}_2$ and $|\boldsymbol{r}_1|^2-|\boldsymbol{r}_2|^2$ multipliers), in which the error actually occurs, and thus, are not affected. The contribution of the coordinate-dependent terms is insignificant for $ka < 0.5$, whereas for larger ka, where the contribution gradually increases, the Q itself becomes too low to be reliably related to the bandwidth.

Further numerical results exemplifying and substantiating the general applicability of a procedure for determining the lower bound on Q outlined in Section V in [1] can be found in [2].

The expressions for the stored energies and the radiated power of arbitrary electric and magnetic currents presented in [1] (Tables I and II) can also be used for computing the Q of electrically small antennas loaded with magneto-dielectric materials, as demonstrated in [3] in the context of a surface integral equation method.

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